

# SAMMY

Goran Arbanas, ORNL

ORNL Collaborators:

Dorothea Wiarda

Luiz C. Leal

Marco Pigni

Mark L. Williams

Vlad Sobes

Andrew M. Holcomb

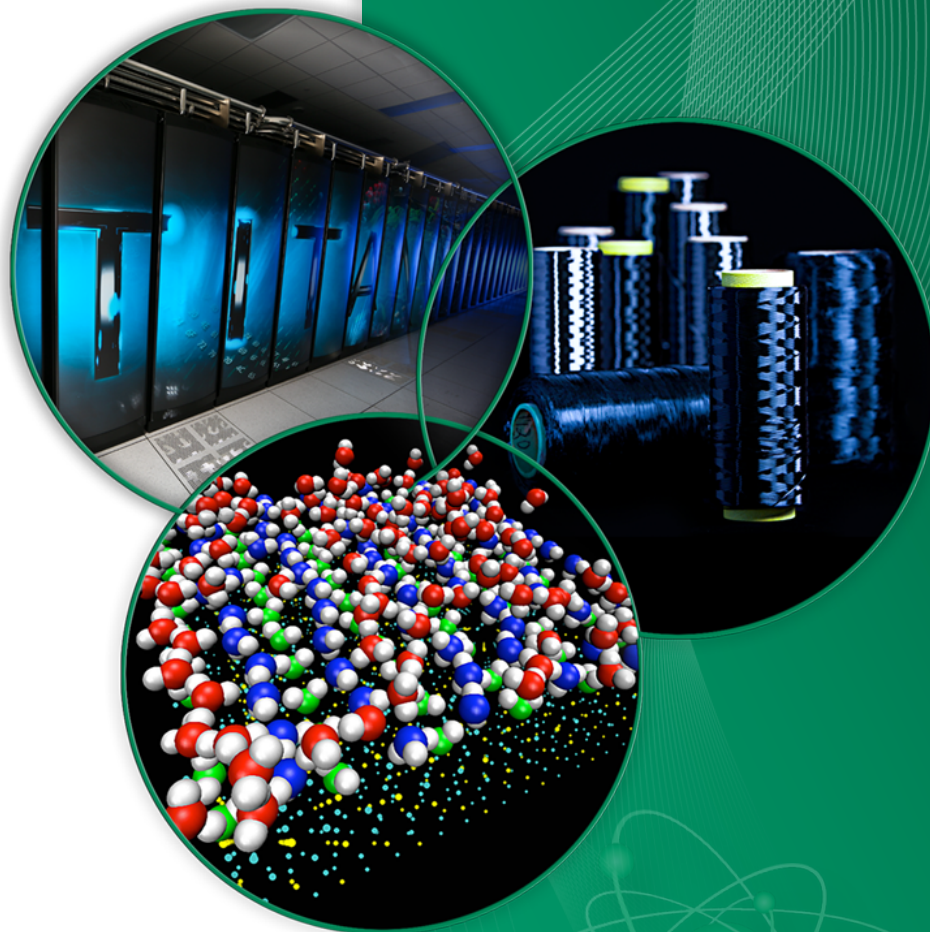
Micheal E. Dunn

RPI Collaborators:

Amanda Youmans,

Kemal Ramic,

Yaron Danon



**Nuclear Data Week, Brookhaven National Laboratory, Nov. 3-7, 2014**

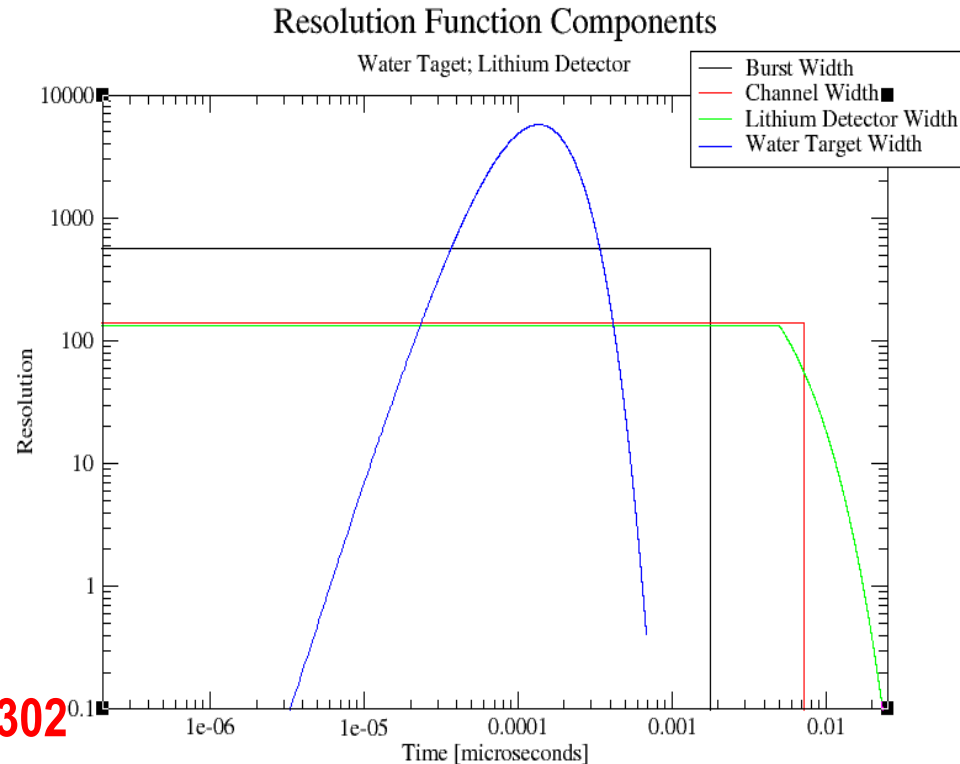
# Overview

- New SAMMY Resolution Functions for RPI neutron detectors
  1. Lithium glass neutron detector array
  2. Liquid scintillator proton recoil neutron detector
- SAMMY Modernization Plan
- S(a,b) project update

# Total SAMMY Resolution Function

- Convolution of four components:

1. Burst width
2. Neutron Sources:
  - Tantalum Target
  - Water Moderator
3. time-of-flight channel width,
4. Neutron Detector:
  - **Lithium Glass Neutron Detector**
  - **Liquid Scintillator Proton Recoil EJ-302**



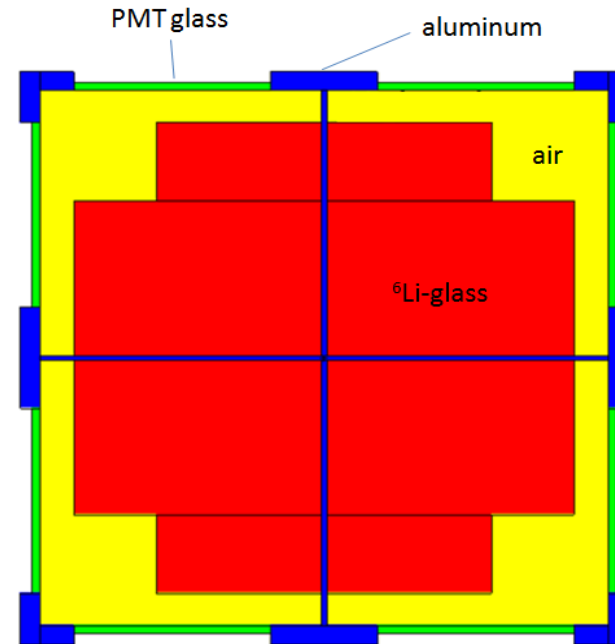
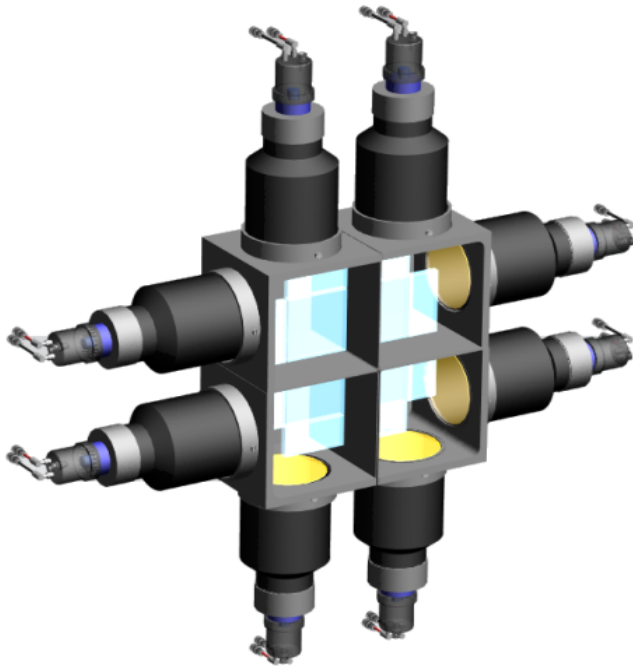
$$\bar{\sigma}(t) = \int I(t-t') \sigma(t') dt'$$

$$E = \frac{1}{2}mv^2 = \frac{1}{2}m \frac{L^2}{t^2}$$

$$I(t-t') = \int I_1(t-t_1) dt_1 \int I_2(t_1-t_2) dt_2 \int I_3(t_2-t_3) dt_3 \int I_4(t_3-t')$$

# $^6\text{Li}$ Glass Detector Array MELINDA

- MCNP simulated the detector and the time distribution of neutrons, tallying  $^6\text{Li}$  (n, $\alpha$ ) reactions



- [1] R. Bahran et al. "A New High Energy Resolution Neutron Transmission Detector System at the Gaertner LINAC Laboratory" *Proceedings of The Tenth International Topical Meeting on Nuclear Applications and Utilization of Accelerators AccApp'11*, Knoxville, TN, 2011.
- [2] A. Youmans et. al. "Using Simulations To Determine The Energy Resolution Function Of Neutron Time-Of-Flight Experiments" *Proceedings of the 2013 Annual ANS Meeting, Washington, D.C., 2013*

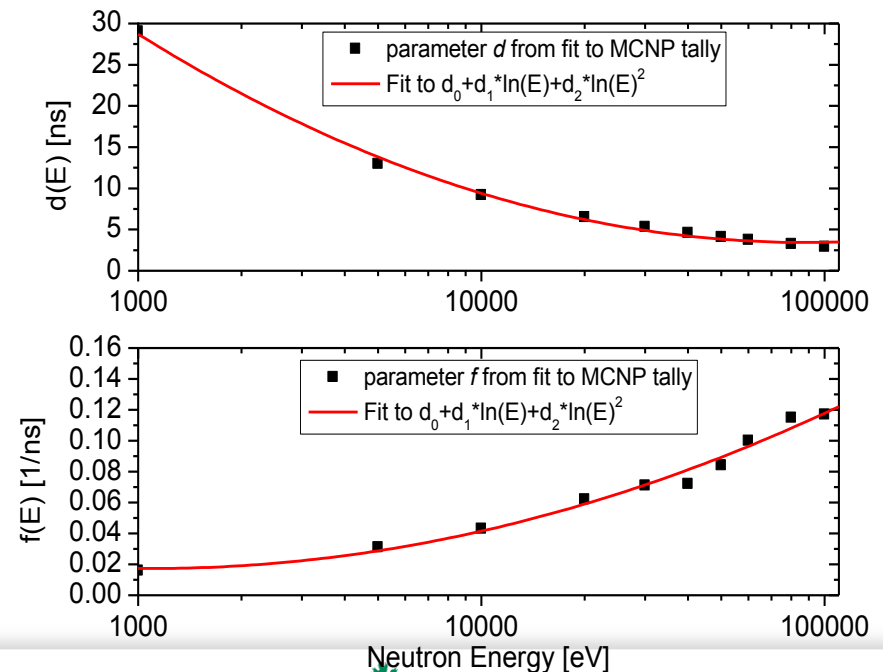
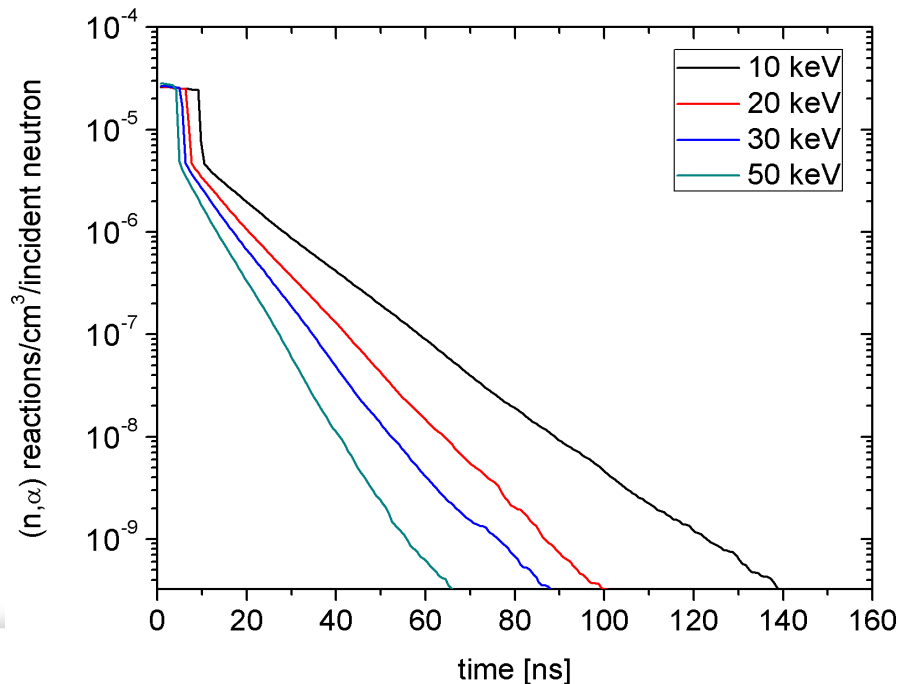


# <sup>6</sup>Li Glass Detector Array MELINDA cont'd.

- From MCNP simulation the RPI resolution function was implemented by making parameters of the SAMMY ORR,  $d$  and  $f$ , energy-dependent:

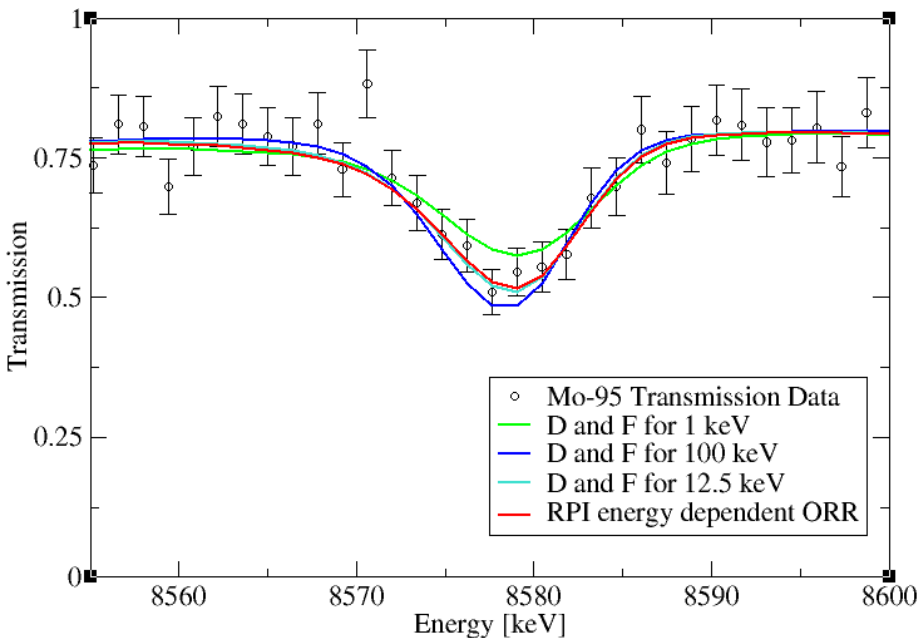
$$I_{4b}(l) = \begin{cases} Dg & \text{for } 0 < t < d \\ D \exp(-f(t-d)) & \text{for } d < t \end{cases}$$

Caveat:  $d_0, d_1, d_2, f_0, f_1, f_2$ , should not be fitted because constant values for  $d$  and  $f$  have been hardwired into SAMMY resolution function.

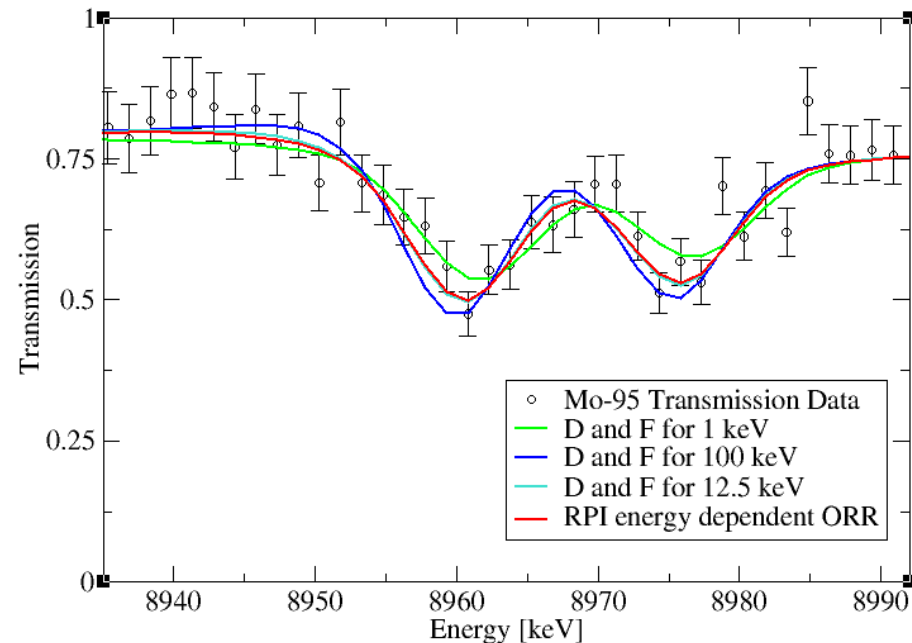


# Results: Li-glass detector

- Sample comparison on RPI Mo-95 transmission data
  - $E$ -dependent vs. constant  $d$  and  $f$
  - We choose three constant values of  $d$  and  $f$  for comparison



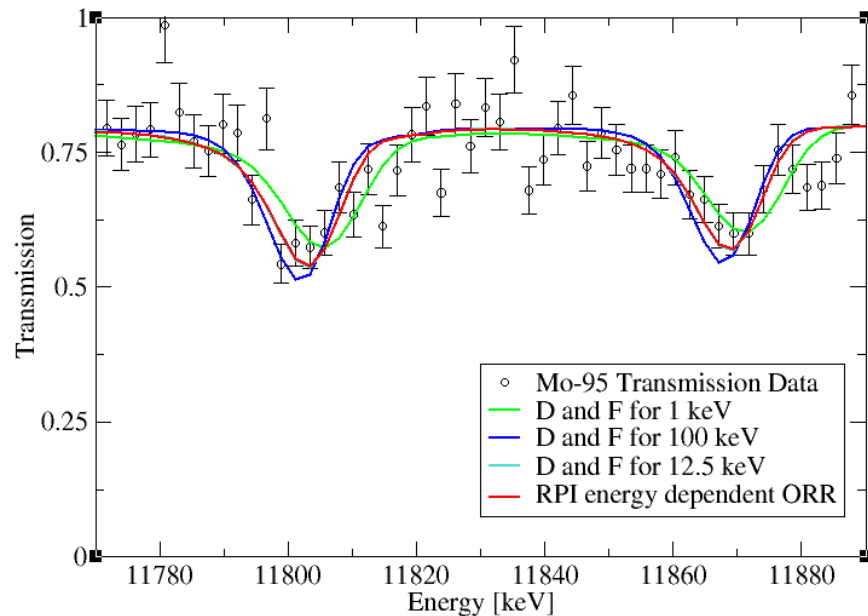
Comparing the original SAMMY detector resolution function to the modified function with Mo-95 data



Comparing the original SAMMY detector resolution function to the modified function with Mo-95 data

# Results: Li-glass resolution function

- E-dependent Li-glass detector resolution function reduces  $\chi^2$

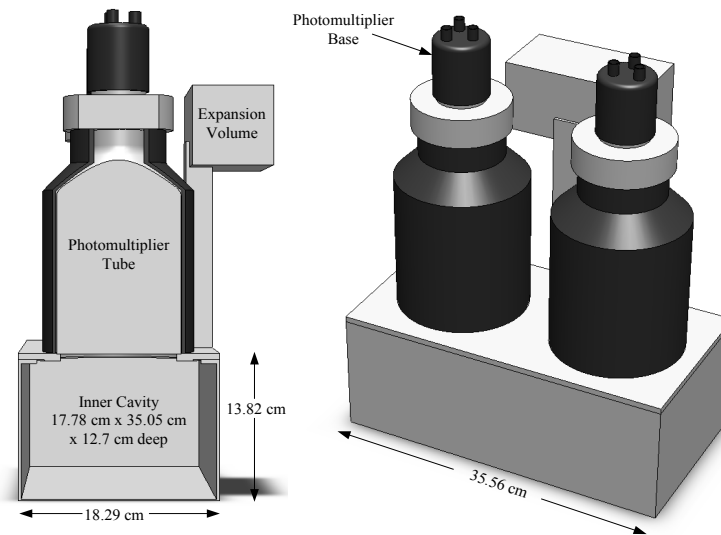


Comparing the original SAMMY detector resolution function to the modified function with Mo-95 data

Trial	Chi-Squared divided by N
$d(E)$ and $f(E)$	1.38422
$d$ , $f$ constant, $E=1$ keV	1.67559
$d$ , $f$ const. $E=12.5$ keV	1.42627
$d$ , $f$ const. $E=100$ keV	1.55281

# EJ-301 Liquid Scintillator

- (MCNP) was used to simulate the detector and the time distribution of neutrons as they react with the detector material, tallying elastic scattering with hydrogen in the proton recoil detector.
- Six modules shown below, each contains liquid  $\text{CH}_{1.212}$ . Neutrons scatter off of hydrogen, ejecting a proton. The proton deposits its energy through charged particle interactions into the scintillation material → pulse



[3] M. Rapp, "Design and construction of a large area detector system and neutron total cross section measurements in the energy range 0.4 to 20 MeV," Ph.D.

dissertation, Mech. Aero. Nucl. Eng. Dept., RPI, Troy, NY, 2011

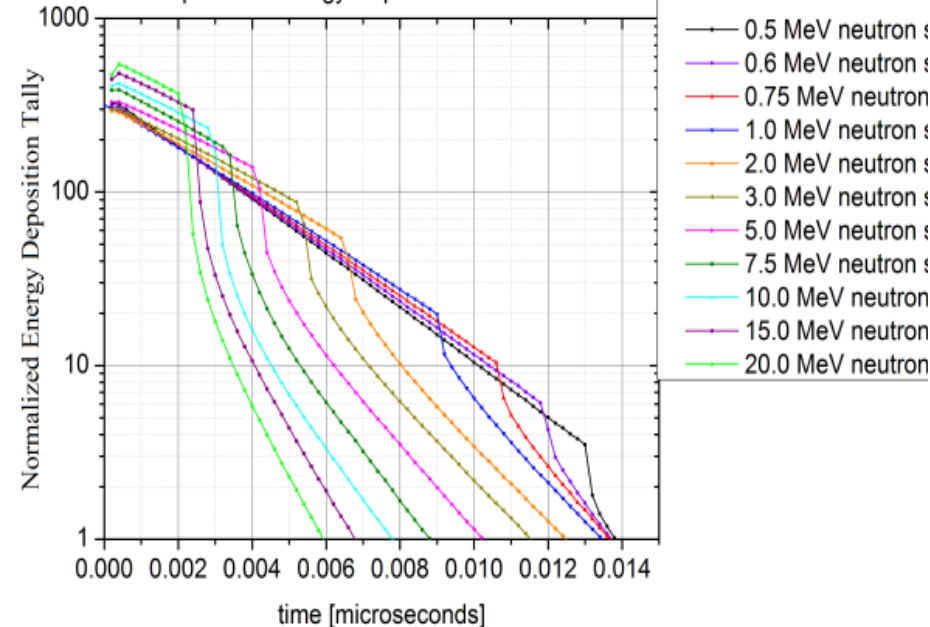
# EJ-301 Liquid Scintillator cont.'d

$$I'_{4a}(l) = \begin{cases} \Delta \exp(-\lambda \sigma l) & \text{for } 0 < l < \delta \\ 0 & \text{otherwise} \end{cases}$$

$$E = \frac{1}{2}mv^2 = \frac{1}{2}m \frac{L^2}{t^2}$$

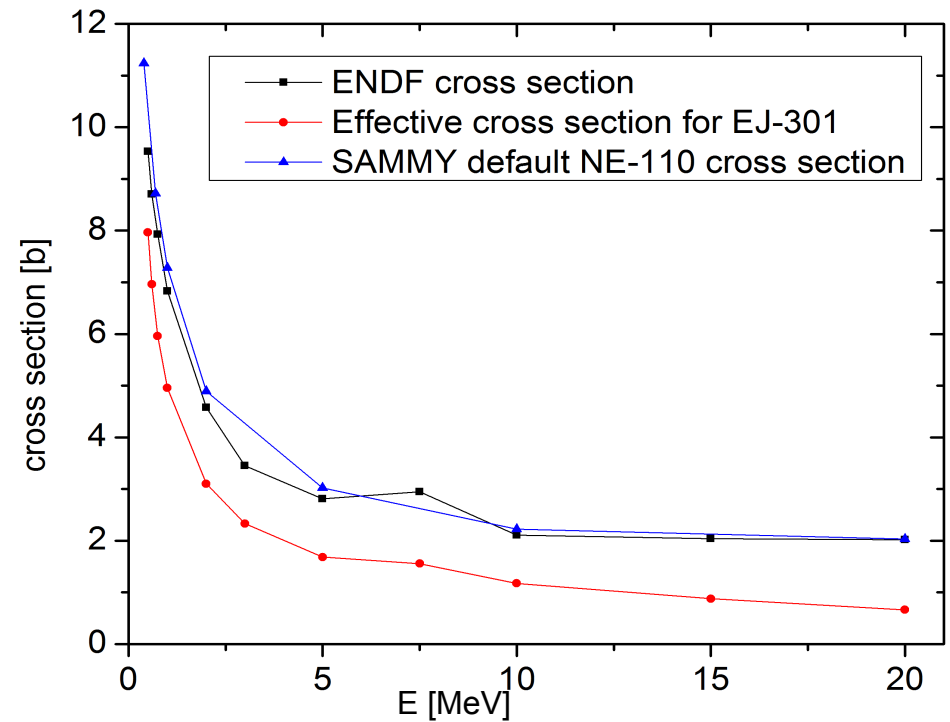
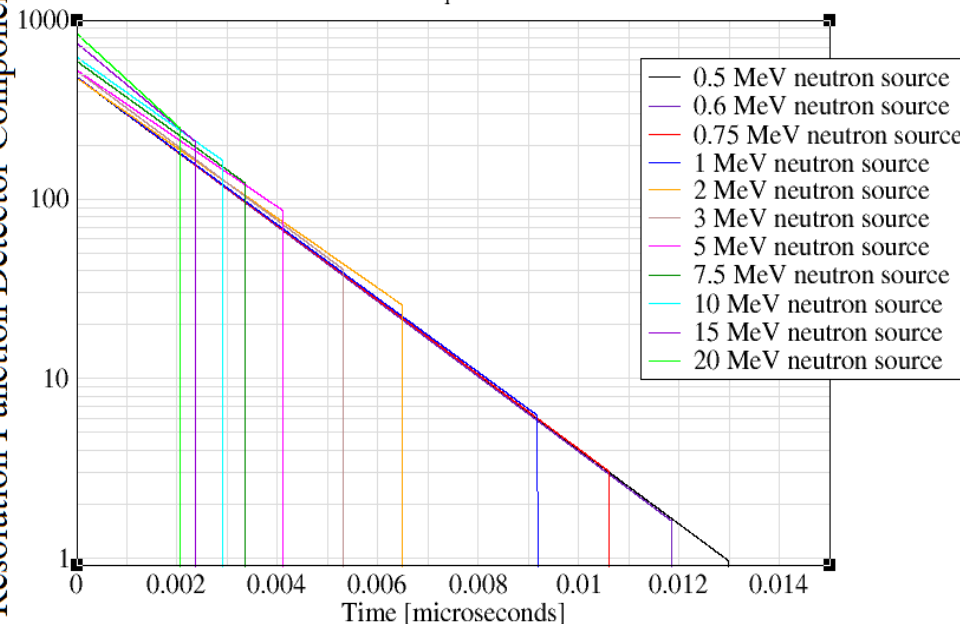
Multiple scattering  
yields tapering below:

MCNP Simulation:  
Time Dependent Energy Deposition in EJ-301 Detector



Resolution Function Detector Component

Default SAMMY Detector Resolution Component  
for NE110 liquid scintillator

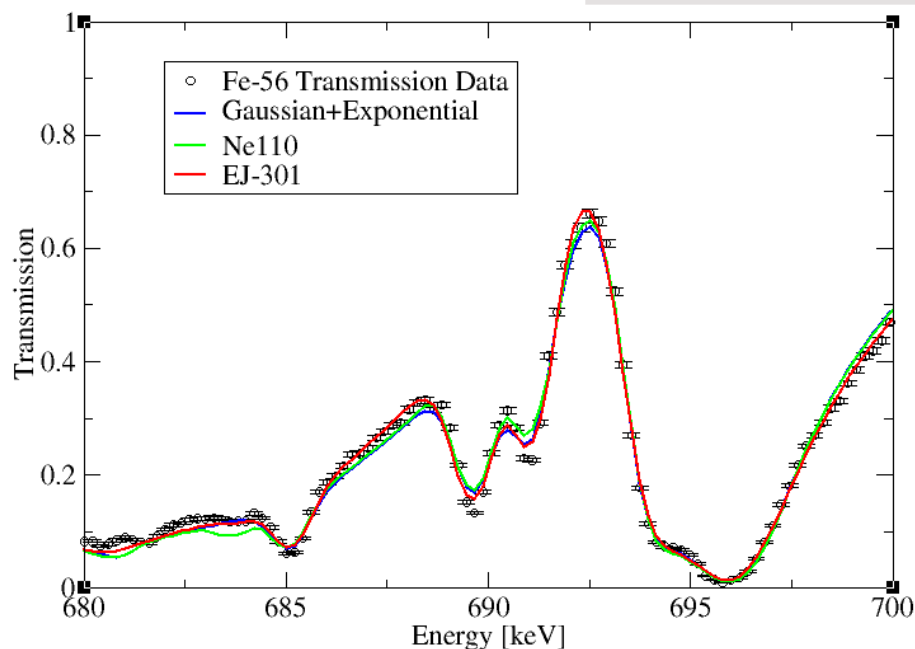


# Results: Liq. Scintil. (EJ-301)

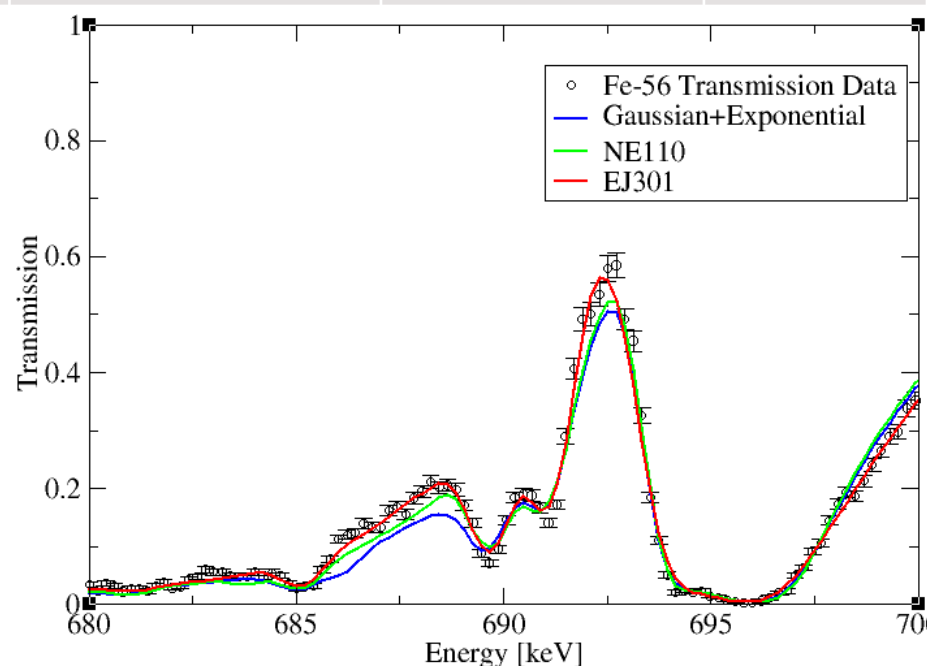
Chi squared divided by N

Resolution Function	Thin Fe56 Sample	Medium Fe56	Thick Fe56
Gauss.+Exp.	13.65	21.42	5.89
Ne-110	12.74	17.31	4.97
EJ-301	6.41	8.26	2.77

- EJ-301 resolution func. reduces  $\chi^2$  (everything else being the same)



Comparison of different detector resolution functions with data from the medium Fe-56 sample at RPI.



Comparison of different detector resolution function with data from the thick Fe-56 sample at RPI.

# SAMMY Modernization Plan

- Use Libraries to simplify programming and parallelization:
  - BLAS, LAPACK, Trillinos (?)
    - Abstractions: e.g. Armadillo library to write matrix product as  $A*B$ ; good for R-matrix
  - Complex numbers, derivatives, etc.
- Implement in C++ (most likely)
  - ORNL SCALE development, build, Q&A framework
- Now is the time to think about connecting to URR, HE, IBEs, etc.
  - Synergies with AMPX
- Separate Bayesian fitting modules from R-matrix computations
  - A generic Bayesian fitting module could also be used for  $S(a,b)$  evaluations
- Feedback from community more than welcome



# S(a,b) methods update

- S(a,b) evaluations via MD appear possible
  - Perform MD simulation
  - For liquids:
    - Compute Intermediate structure function (simpler than dyn. struct. funct.)
    - From it, compute dynamical structure function S(k,w)
  - For solids: one could use LAMMPS/PHANA to get phonon spectrum
  - ORNL Spallation Neutron Source data: S(k,w) and its covariance
    - compare and iterate, GLLS will yield interaction parameters and their covariance
  - Compute sensitivities from an adjoint MD simul, not numerical variations
- Applied for and received NERSC Startup Allocation: only 50K but:
  - All MD and DFT codes (including *ab initio* ones) available at NERSC
  - Will enable a prototype evaluation that could lead to other applications
- Promising collaborations here and with ORNL's MD experts
  - Will apply for ORNL's OLCF Director's Discretionary supercomp alloc.

Third-party  
codes failing  
for water MD

This  
works

# Additional slides

# Overview

- Goal: Develop methods for generating covariance data of  $S(a,b)$ 
  - In anticipation of the RPI data for  $H_2O$  and  $SiO_2$  (Si), polyethylene; SNS
- General approach:
  1. Define model  $T$  as a function of some parameters  $P$
  2. Fit model  $T$  to data to obtain covariance matrix  $M$  of parameters  $P$
  3. Construct covariance of  $S(a,b)$  from  $M$  as  $C = S M S$ 
    - Where  $S = \partial T(P) / \partial P$  is the sensitivity matrix of model  $T$
- Various models  $T$  being considered:
  - Addressing various levels of physical models
    - Various approximations: from DFT to MD to free gas,
    - Some assumptions revisited (e.g. form of the scattering w.f.)
  - *Ab initio* models more computationally intensive
- MADNESS Computational framework (R&D 100, SciDAC)
  - May provide novel computational approaches and insights

# General consideration

- The form of the scattering function assumed
  - Plane wave + scattered wave
  - Can we estimate the magnitude of ensuing error? (assumed to be small)
- The low-energy ( $< 5$  eV) neutron-nucleus scattering determined by the scattering length
- First-order Born approximation
  - Does not require a delta-function interaction;
    - Woods-Saxon is more realistic
    - But delta-function is likely a good approximation since  $\text{size}(\text{nuclei}) \ll \text{size}(\text{atoms})$
- This leads to expressions for  $S(k,w)$  or  $S(a,b)$  in terms of pair-wise correlation functions
  - Delineates neutron-scattering from material properties PCF
  - $\rightarrow$  PCF could be computed by e.g. MD
- What is an efficient path to  $S(a,b)$  covariance?

# 0. Single molecule scattering

- The motion of hydrogen atoms in water in terms of the H<sub>2</sub>O molecule as the basic dynamical unit.
  - Mark Nelkin, Phys. Rev. 119, **741** (1960)
- We could fit parameters of the  $S(a,b)$  expression to data
  - Parameters of water vibrations, torsional oscillations, etc.
  - From covariance/correlation of fitted parameters get covariance of  $S(a,b)$
- We could also fit *neutron-nucleus* potential parameters

# 1. Molecular Dynamics (MD)

- MD Method:
  - assumes a parameterized interaction potential (e.g. Lennard-Jones)
  - Uses Newton's equation to evolve the many-body system
  - Then computes PCFs to get  $S(a,b)$  and structure factors
  - Faster but less accurate than more *ab initio* methods like DFT
  - Interaction parameters fitted to (usually) structure factors  $S(k)$
  - The  $S(k)$  is the 0-th moment of  $S(q,w)$
  - But parameter uncertainty or covariance is generally not provided
- Several established MD codes available
  - GROMACS, LAMMPS, NAMD
  - [http://en.wikipedia.org/wiki/List\\_of\\_software\\_for\\_molecular\\_mechanics\\_modeling](http://en.wikipedia.org/wiki/List_of_software_for_molecular_mechanics_modeling)
  - Output molecular/atomic trajectories in time

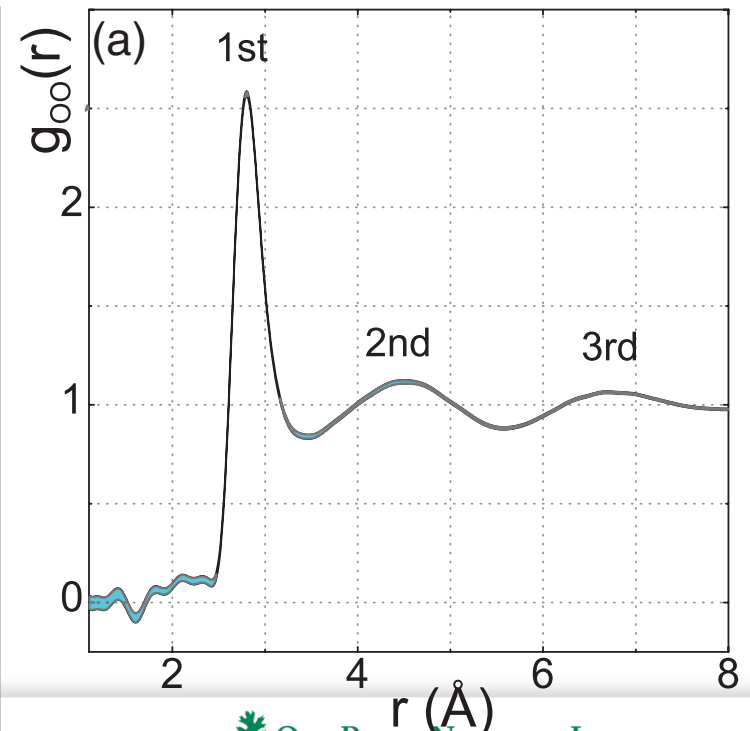
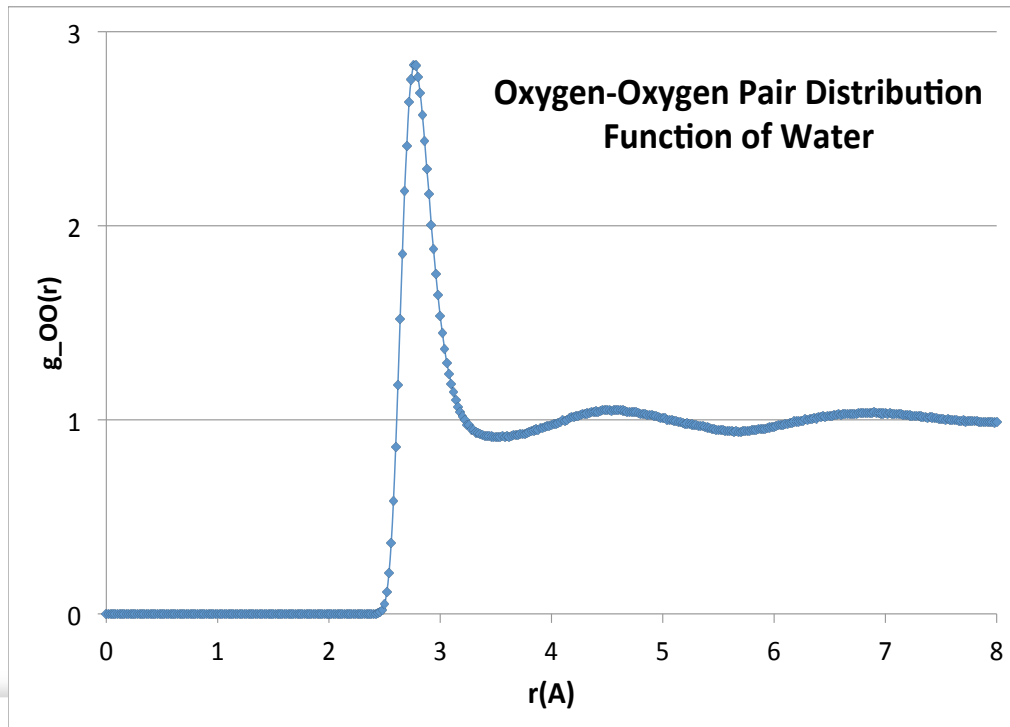
# 1. Molecular Dynamics (MD) cont'd.

- Trajectories processed to compute  $G(r,t)$ ,  $I(k,t)$ ,  $S(k,w)$ 
  - Using dynsf: can take days to complete, prone to crashing
- Multi-atomic molecules like  $H_2O$  require computations  $S(k,w)$  for all possible different pairs:
  - O-O, H-H, O-H
  - The total  $S(k,w)$  is the weighted average of the the partial  $S(k,w)$ 
    - Weights are: ( scattering lengths ) \* ( elemental fractions )
- Currently computing pairwise  $S(k,w)$  using dynsf
  - may need to develop a faster version
  - Cf. LAMMPS MD simulation takes a couple of hours on 32 processors
- Computed polyethylene, and  $SiO_2$  phonon spectrum using LAMMPS
  - Also PE phonon spectrum has been extracted from SEQUOIA data
  - Data fitting and covariances are on the horizon



# Gromacs Example Run Water

- O-O radial distribution function  $g(r)$ 
  - Classical pair correlation function  $G(r,t)$  would give  $S(k,w)$
  - [https://github.com/cfinch/GROMACS\\_Examples/tree/master/water](https://github.com/cfinch/GROMACS_Examples/tree/master/water)
  - Must use third party code e.g. dynsf to compute  $G(r,t)$  from the trajectories



## 2. Density Functional Theory

- Codes:
  - VASP, QuantumEspresso, BigDFT, MADNESS (more later)
  - [http://en.wikipedia.org/wiki/List\\_of\\_quantum\\_chemistry\\_and\\_solid\\_state\\_physics\\_software](http://en.wikipedia.org/wiki/List_of_quantum_chemistry_and_solid_state_physics_software)
- Water is still not fully understood
  - Discrepancies between data and theory persist
    - structure factor, pairwise correlation function
  - Several recent Ph.D. thesis on water (active field)
- ~1,000's of atoms/molecules
- More accurate than MD
- DFT computations are used to fit MD interaction parameters
- A candidate framework for  $S(a,b)$  covariances

### 3. Total Montel Carlo

- By Koning and Rochman
- Could compute *covariance* of existing  $S(a,b)$  evaluations
  - For which covariance data does not exist
  - This method yields covariance by averaging large random ensembles
- Could be compared to new evaluations

## 4. SNS SEQUOIA Data reduction

Horace platform:

- One of the two preferred methods for data reduction and analysis
  - at the SNS SEQUOIA
  - double diff. (angle, energy) cross section is reduced to  $S(\mathbf{q}, \omega)$
  - Built-in function for parameter fitting to reduced data  $S(\mathbf{q}, \omega)$ 
    - Returns model parameters, their uncertainties, and the correlation matrix
    - One can compute the model covariance matrix of the  $S(\mathbf{q}, \omega)$
- An established platform
  - Works with MATLAB out of the box
  - Used in multitude of advanced papers
  - Downloadable from <http://horace.isis.rl.ac.uk>

Also DAVE: <http://www.ncnr.nist.gov/dave/>

# Double-Differential Cross Section

- Ouisloumen and Sanchez (1991) cf. Rothenstein/Dagan/Becker
- Arbanas *et al.* M&C 2011, NCSC2 2012, PHYSOR 2012

$$\sigma_s^T(E \rightarrow E', \mu_{\text{lab}}) = \sum_{n \geq 0} (n + \frac{1}{2}) \sigma_{sn}^T(E \rightarrow E') P_n(\mu_{\text{lab}})$$

$$\sigma_{sn}^T(E \rightarrow E') = \frac{\beta^{5/2}}{4E} e^{E/kT} \int_0^\infty t \sigma_s^{\text{tab}}(E''(t)) e^{-t^2/A} \psi_n(t) dt,$$

$$\sigma_{sn}(\mathbf{x}) = \hat{O}_n(\mathbf{x})[\sigma_s^{\text{tab}}], \quad \mathbf{x} \equiv (T, E, E'),$$

$$M_{nn'}(\mathbf{x}, \mathbf{y}) \equiv \langle \delta \sigma_{sn}(\mathbf{x}) \delta \sigma_{sn'}(\mathbf{y}) \rangle = \hat{O}_n(\mathbf{x}) \hat{O}_{n'}(\mathbf{y}) [V_s]$$

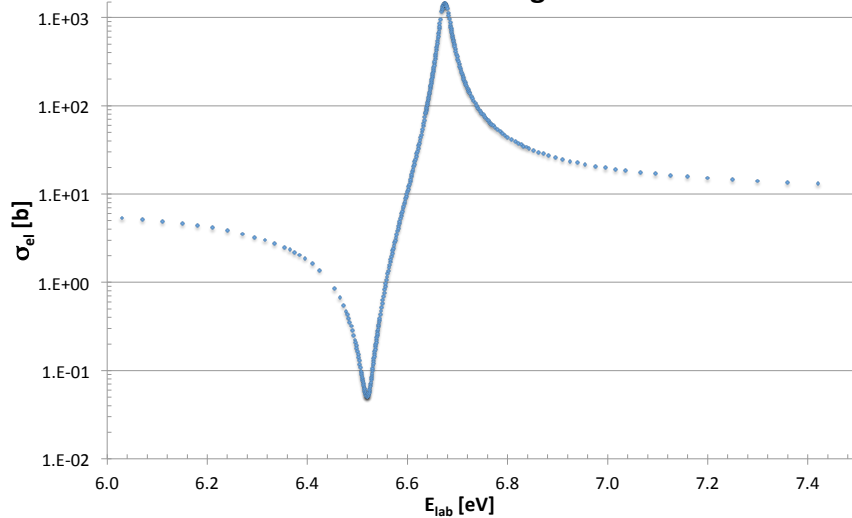
$$V_s(E, E') \equiv \langle \delta \sigma_s^{\text{tab}}(E) \delta \sigma_s^{\text{tab}}(E') \rangle$$

$$\langle \delta \sigma_s(\mathbf{x}, \mu_{\text{lab}}) \delta \sigma_s(\mathbf{y}, \mu'_{\text{lab}}) \rangle = \sum_{n, n'} (n + \frac{1}{2})(n' + \frac{1}{2}) \times P_n(\mu_{\text{lab}}) P_{n'}(\mu'_{\text{lab}}) M_{nn'}(\mathbf{x}, \mathbf{y})$$

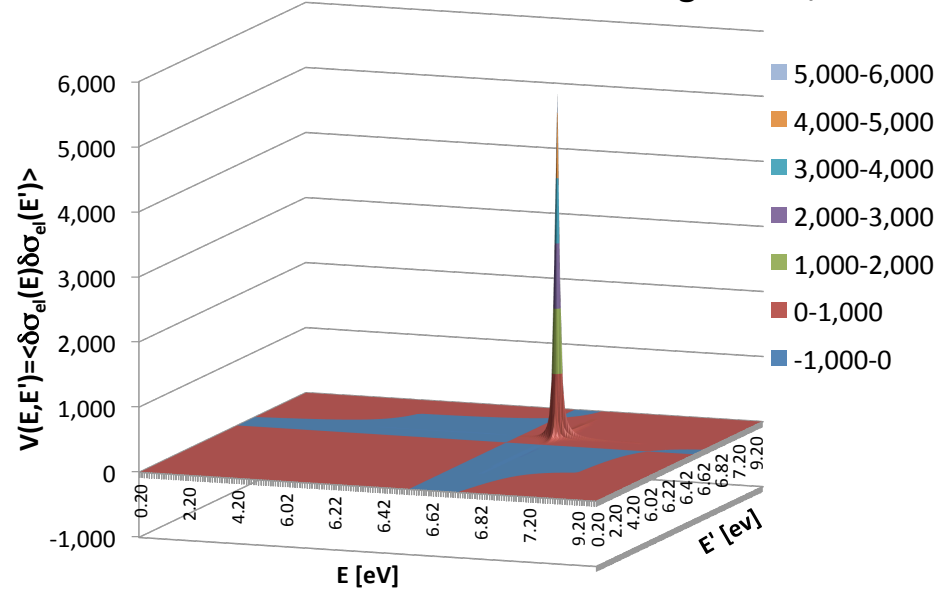
- The above could be used to connect to the thermal S(a,b)

# Doppler Broadened Legendre Moments

$^{238}\text{U}$  Neutron Elastic Scattering Cross Section

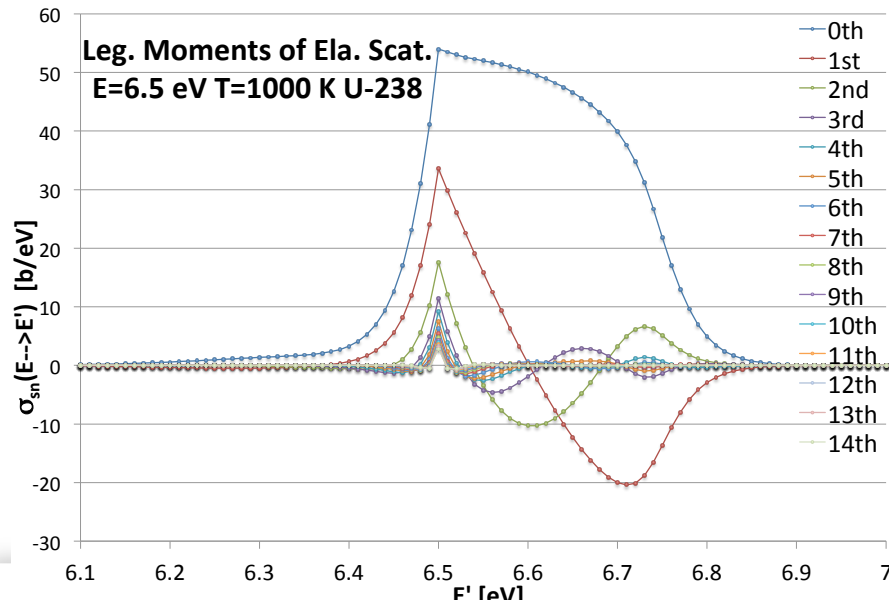


Covariance Matrix of Elastic Scattering  $n+^{238}\text{U}$ ,  $T=0$  K



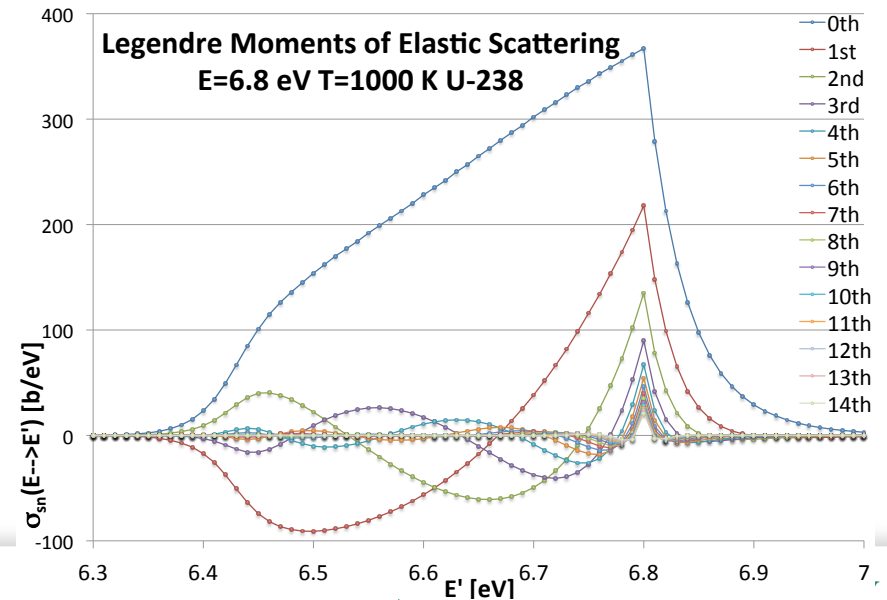
Leg. Moments of Ela. Scat.

$E=6.5$  eV  $T=1000$  K  $\text{U-238}$



Legendre Moments of Elastic Scattering

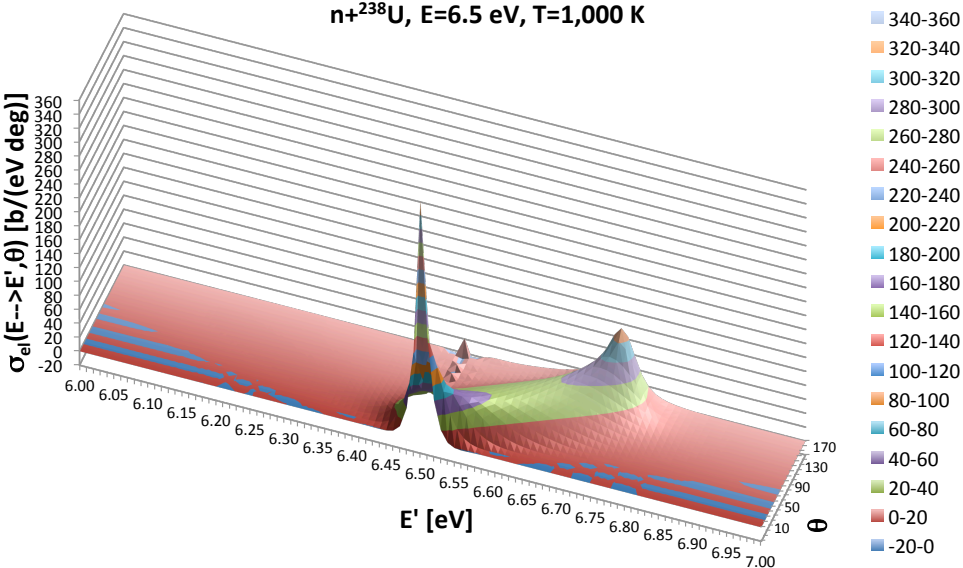
$E=6.8$  eV  $T=1000$  K  $\text{U-238}$



# Double-Differential Cross Section

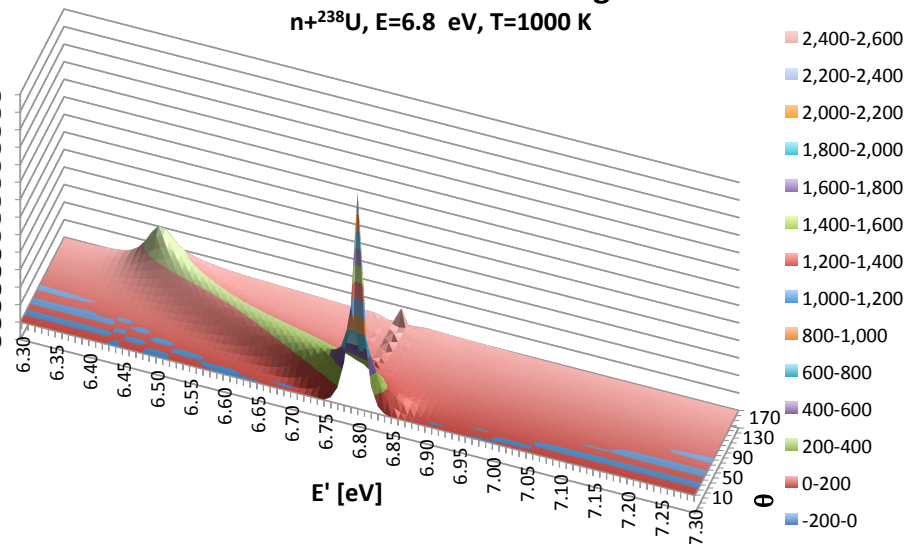
## Double-Differential Elastic Scattering Cross Section

$n+^{238}\text{U}$ ,  $E=6.5$  eV,  $T=1,000$  K



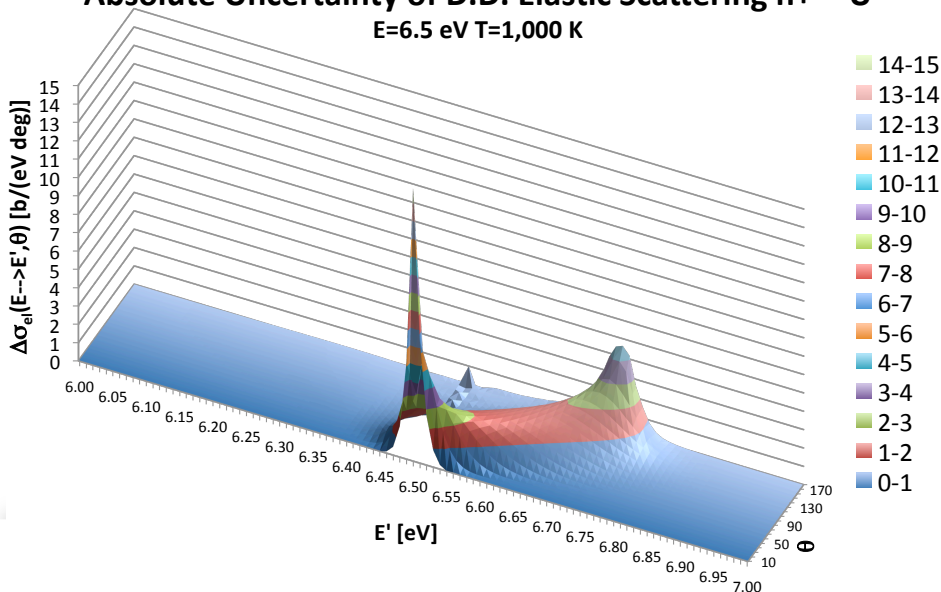
## Double-Differential Elastic Scattering Cross Section

$n+^{238}\text{U}$ ,  $E=6.8$  eV,  $T=1000$  K



## Absolute Uncertainty of D.D. Elastic Scattering $n+^{238}\text{U}$

$E=6.5$  eV  $T=1,000$  K



## Absolute Uncertainty of D.D. Elastic Scattering $n+^{238}\text{U}$

$E=6.8$  T=1000 K

